

Modern Physics: Chapter 6 Homework

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1. (book #2) Show that

$$\Theta_{20}(\theta) = \frac{\sqrt{10}}{4}(3 \cos^2 \theta - 1)$$

is a solution of Equation (6.13) and that it is normalized.

We need to show that

$$\frac{1}{\sin \theta} \frac{d}{d\theta} \left(\sin \theta \frac{d\Theta}{d\theta} \right) + \left[l(l+1) - \frac{m_l^2}{\sin^2 \theta} \right] \Theta = 0 \quad (1)$$

for $l = 2$, $m_l = 0$, and for all values of θ .

We have $\frac{d\Theta}{d\theta} = \frac{-3\sqrt{10}}{2} \cos \theta \sin \theta$, so that the expression on the left hand side of (1) above is

$$\begin{aligned} & \frac{1}{\sin \theta} \frac{d}{d\theta} \left(\frac{-3\sqrt{10}}{2} \cos \theta \sin^2 \theta \right) + \left[6 \frac{\sqrt{10}}{4} (3 \cos^2 \theta - 1) \right] \\ &= \frac{1}{\sin \theta} \left(\frac{-3\sqrt{10}}{2} \right) (2 \cos^2 \theta \sin \theta - \sin^3 \theta) + \frac{3\sqrt{10}}{2} (3 \cos^2 \theta - 1) \\ &= \frac{-3\sqrt{10}}{2} (2 \cos^2 \theta - \sin^2 \theta) + \frac{3\sqrt{10}}{2} (3 \cos^2 \theta - 1) \\ &= 0 \quad \text{since } \sin^2 \theta + \cos^2 \theta = 1 \implies \sin^2 \theta = 1 - \cos^2 \theta \end{aligned}$$

To show that Θ is normalized, we will show that $\int_0^\pi |\Theta|^2 \sin \theta d\theta = 1$.

We have,

$$\begin{aligned} \int_0^\pi |\Theta|^2 \sin \theta d\theta &= \int_0^\pi \frac{10}{16} (3 \cos^2 \theta - 1)^2 \sin \theta d\theta \\ &= -\frac{5}{8} \int_1^{-1} (3u^2 - 1)^2 du \quad \text{with } u = \cos \theta, \quad du = -\sin \theta d\theta \\ &= -\frac{5}{8} \int_1^{-1} (9u^4 - 6u^2 + 1) du \\ &= -\frac{5}{8} \left[\frac{9u^5}{5} - \frac{6u^3}{3} + u \right]_1^{-1} \\ &= -\frac{5}{8} \left(-\frac{9}{5} + 2 - 1 - \frac{9}{5} + 2 - 1 \right) \\ &= -\frac{5}{8} \left(-\frac{8}{5} \right) = 1 \end{aligned}$$

Q.E.D.

We will use the following Lemma in problems #2 and #8:

Lemma 1. $\int_0^\infty r^k e^{-pr} dr = \left(\frac{1}{p}\right)^{k+1} k!$ for $k \in N, p > 0$

Let $u = r^k$, $dv = e^{-pr} dr$ and hence $du = kr^{k-1} dr$, $v = \frac{-1}{p}e^{-pr}$. Then,

$$\begin{aligned} \int_0^\infty r^k e^{-pr} dr &= \int_0^\infty u dv \\ &= \left[r^k \left(-\frac{1}{p} e^{-pr} \right) \right]_0^\infty + \int_0^\infty \frac{1}{p} e^{-pr} \cdot kr^{k-1} dr \\ &= \frac{k}{p} \int_0^\infty r^{k-1} e^{-pr} dr \\ &= \left(\frac{1}{p} \right)^k k! \int_0^\infty e^{-pr} dr \\ &= - \left(\frac{1}{p} \right)^{k+1} k! [e^{-pr}]_0^\infty \\ &= \left(\frac{1}{p} \right)^{k+1} k! \end{aligned}$$

Q.E.D.

2. (book #4) Show that

$$R_{21}(r) = \frac{1}{2\sqrt{6}a_0^{3/2}} \frac{r}{a_0} e^{-r/2a_0}$$

is a solution of Equation (6.14) and that it is normalized.

We need to show that

$$\frac{1}{r^2} \frac{d}{dr} \left(r^2 \frac{dR}{dr} \right) + \left[\frac{2m}{\hbar^2} \left(\frac{e^2}{4\pi\epsilon_0 r} + E \right) - \frac{l(l+1)}{r^2} \right] R = 0$$

for $n = 2, l = 1$, and for all values of r .

We have

$$\begin{aligned} \frac{dR}{dr} &= \frac{1}{2\sqrt{6}a_0^{5/2}} \frac{d}{dr} \left(r e^{-r/2a_0} \right) \\ &= \frac{1}{2\sqrt{6}a_0^{5/2}} \left(\frac{-r}{2a_0} e^{-r/2a_0} + e^{-r/2a_0} \right) \\ &= \frac{e^{-r/2a_0}}{2\sqrt{6}a_0^{5/2}} \left(\frac{-r}{2a_0} + 1 \right) \end{aligned}$$

Then

$$\begin{aligned}
\frac{d}{dr} \left(r^2 \frac{dR}{dr} \right) &= \frac{d}{dr} \left[\frac{r^2 e^{-r/2a_0}}{2\sqrt{6}a_0^{5/2}} \left(\frac{-r}{2a_0} + 1 \right) \right] \\
&= \frac{1}{2\sqrt{6}a_0^{5/2}} \frac{d}{dr} \left[r^2 e^{-r/2a_0} \left(\frac{-r}{2a_0} + 1 \right) \right] \\
&= \frac{1}{2\sqrt{6}a_0^{5/2}} \left[\frac{-r^2 e^{-r/2a_0}}{2a_0} + \left(\frac{-r}{2a_0} + 1 \right) \left(\frac{-r^2 e^{-r/2a_0}}{2a_0} + 2r e^{-r/2a_0} \right) \right] \\
&= \frac{r e^{-r/2a_0}}{2\sqrt{6}a_0^{5/2}} \left[\frac{-r}{2a_0} + \left(\frac{-r}{2a_0} + 1 \right) \left(\frac{-r}{2a_0} + 2 \right) \right]
\end{aligned}$$

so that

$$\begin{aligned}
\frac{1}{r^2} \frac{d}{dr} \left(r^2 \frac{dR}{dr} \right) &= \frac{e^{-r/2a_0}}{2\sqrt{6}a_0^{5/2}} \left[\frac{-1}{2a_0} + \frac{r}{4a_0^2} - \frac{1}{a_0} - \frac{1}{2a_0} + \frac{2}{r} \right] \\
&= \frac{e^{-r/2a_0}}{2\sqrt{6}a_0^{5/2}} \left(\frac{-2}{a_0} + \frac{r}{4a_0^2} + \frac{2}{r} \right)
\end{aligned}$$

Meanwhile,

$$\left[\frac{2m}{\hbar^2} \left(\frac{e^2}{4\pi\epsilon_0 r} + E \right) - \frac{l(l+1)}{r^2} \right] R = \left[\frac{2m}{\hbar^2} \left(\frac{e^2}{4\pi\epsilon_0 r} + E \right) - \frac{2}{r^2} \right] \frac{r e^{-r/2a_0}}{2\sqrt{6}a_0^{5/2}}$$

so that it suffices to show that

$$-\frac{2}{a_0} + \frac{r}{4a_0^2} + \frac{2}{r} + \frac{2mr}{\hbar^2} \left(\frac{e^2}{4\pi\epsilon_0 r} + E \right) - \frac{2}{r} = 0$$

On the left hand side we have

$$\frac{-2}{a_0} + \frac{r}{4a_0^2} + \frac{2mre^2}{4\hbar^2\pi\epsilon_0 r} + \frac{2mrE}{\hbar^2}$$

Then substituting for $a_0 = \frac{4\pi\hbar^2\epsilon_0}{me^2}$ we have

$$\begin{aligned}
\frac{-2me^2}{4\pi\hbar^2\epsilon_0} + \frac{rm^2e^4}{4 \cdot 16\pi^2\hbar^4\epsilon_0^2} + \frac{2me^2}{4\hbar^2\pi\epsilon_0} + \frac{2mrE}{\hbar^2} &= \frac{rm^2e^4}{64\pi^2\hbar^4\epsilon_0^2} + \frac{2mrE}{\hbar^2} \\
&= \frac{2rm}{\hbar^2} \left(\frac{me^4}{32\pi^2\hbar^2\epsilon_0^2} \left(\frac{1}{2^2} \right) + E \right)
\end{aligned}$$

which is equal to zero precisely when

$$E = -\frac{me^4}{32\pi^2\hbar^2\epsilon_0^2} \left(\frac{1}{n^2} \right)$$

which corresponds to the energy levels of the hydrogen atom. Therefore R is a solution of Equation (6.14).

To show that R is normalized, we will show that

$$\int_0^\infty r^2 |R|^2 dr = 1$$

We have

$$\begin{aligned} \int_0^\infty r^2 |R|^2 dr &= \int_0^\infty \frac{r^4}{24a_0^5} e^{-r/a_0} dr \\ &= \frac{1}{24a_0^5} \int_0^\infty r^4 e^{-r/a_0} dr \\ &= \frac{1}{24a_0^5} (a_0)^5 4! \quad \text{by Lemma 1} \\ &= 1 \end{aligned}$$

Q.E.D.

3. (book #9) Under what circumstances, if any, is L_z equal to L ?

Only when $L = 0$ can L_z equal L . Otherwise $|L| > |L_z|$.

4. (Book #11) What are the possible values of the magnetic quantum number m_l of an atomic electron whose orbital quantum number is $l = 4$?

Since $l \geq |m_l|$, we have $m_l \in \{-4, -3, -2, -1, 0, 1, 2, 3, 4\}$.

5. (Book #12) List the sets of quantum numbers possible for an $n = 4$ hydrogen atom.

We may have $l \in \{0, 1, 2, 3\}$. Then:

When $l = 0$, $m_l = 0$.

When $l = 1$, $m_l \in \{-1, 0, 1\}$.

When $l = 2$, $m_l \in \{-2, -1, 0, 1, 2\}$.

When $l = 3$, $m_l \in \{-3, -2, -1, 0, 1, 2, 3\}$.

These comprise all quantum numbers possible for an $n = 4$ hydrogen atom.

6. (Book #17) Find the most probable value of r for a 3d electron in a hydrogen atom.

For a 3d electron, we have

$$R_{32}(r) = \frac{4r^2 e^{-r/3a_0}}{81\sqrt{30}a_0^{7/2}}$$

Then the probability of finding the electron at radius r is given by

$$P(r) = r^2 |R|^2 = \frac{4^2 r^6 e^{-2r/3a_0}}{81^2 \cdot 30a_0^7}$$

We need to find r for which $P(r)$ is at maximum. Hence we have

$$\begin{aligned} \frac{d}{dr}P(r) &= \frac{4^2}{81^2 \cdot 30a_0^7} \frac{d}{dr} \left(r^6 e^{-2r/3a_0} \right) = 0 \\ \implies \frac{-2r^6 e^{-2r/3a_0}}{3a_0} + 6r^5 e^{-2r/3a_0} &= 0 \end{aligned}$$

But $r = 0$ is a minimum since $P(0) = 0$. Therefore

$$\begin{aligned} \frac{-2r}{3a_0} + 6 &= 0 \\ \implies 6 &= \frac{2}{3a_0}r \\ \implies r &= 9a_0 \end{aligned}$$

Therefore $\boxed{9a_0}$ is the most probable radius for a 3d electron in a hydrogen atom.

7. (Book #19) How much more likely is the electron in a ground-state hydrogen atom to be at the distance a_0 from the nucleus than at the distance $2a_0$?

For an electron in the ground state,

$$R = \frac{2}{a_0^{3/2}} e^{-r/a_0}$$

Then letting P_1 be the probability that the electron will be at $r_1 = a_0$ and letting P_2 be the probability that the electron will be at $r_2 = 2a_0$, we need to find P_1/P_2 .

We have

$$\frac{P_1}{P_2} = \frac{r_1^2 |R_1|^2}{r_2^2 |R_2|^2} = \frac{(a_0)^2 e^{-2a_0/a_0}}{(2a_0)^2 e^{-4a_0/a_0}} = \frac{1e^{-2}}{4e^{-4}} = \frac{1}{4}e^2 = 1.85$$

Therefore it is $\boxed{1.85}$ times as likely that the electron will be found at radius a_0 than at $2a_0$.

8. (Book #20) Verify that the average value of r for a $1s$ electron in a hydrogen atom is $1.5a_0$ by calculating the expectation value $\langle r \rangle = \int r|\psi|^2 dV$.

We have

$$\begin{aligned}\langle r \rangle &= \int_0^\infty r |R_{10}|^2 r^2 dr \\ &= \int_0^\infty r^3 \left(\frac{2}{a_0^{3/2}} e^{-r/a_0} \right)^2 dr \\ &= \frac{4}{a_0^3} \int_0^\infty r^3 e^{-2r/a_0} dr \\ &= \frac{4}{a_0^3} \left(\frac{a_0}{2} \right)^4 3! \quad \text{by Lemma 1} \\ &= \frac{4a_0}{2^4} (3 \cdot 2) \\ &= \frac{3}{2} a_0\end{aligned}$$

Therefore the average value of r for a $1s$ electron in the hydrogen atom is $1.5a_0$.

Q.E.D.

9. (Book #24) A hydrogen atom is in the $4p$ state. To what state or states can it go by radiating a photon in an allowed transition?

By the selection rule $\Delta l = \pm 1$, the electron beginning in the $n = 4, l = 1$ state may go to any of the following states by radiating a photon: $\boxed{3s, 3d, 2s, 1s}$.